

Fig. 1. The IRC on a two-dimensional model potential surface (--- valley paths and displacement lines, see text)

and change the potential surface by mutual displacements of the two rigid minima-hollows along the direction of the eigenvector of the smallest eigenvalue (that means we only change level lines upwards from the (---)line) we will observe a change in the relative position of the saddle to those of the both minima. Hence, it results a change of the IRC also for the fixed part of the surface. The knowledge of the local curvatures of the potential in a minimum position does not generally allow a conclusion in which direction the lowest next saddle can be found.

It should be denoted that in Fig. 1 no branching of the (---)valley exists. So we cannot follow one of the branches from the minimum to the saddle obtaining the IRC. Consequently, the idea of the optimal ascent path [6] fails here.

(iii) The solutions of the gradient equations (2, 8, or 23) cross the equipotential surfaces rectangular in every regular point [25, 26]. In a potential hollow the equipotential surfaces are convex, hence the potential has a minimal value on the tangential plane in any point x where this plane touches the potential function (see Fig. 2). The touching point be x_1 . - This right insight nevertheless led to considerable trouble, because it suggested a valley-path character ("minimum energy path") of every solution curve $x(t)$. Of course, rectangular to the direction of the descent $-\nabla U$ we have an ascent to both sides. But besides the generally

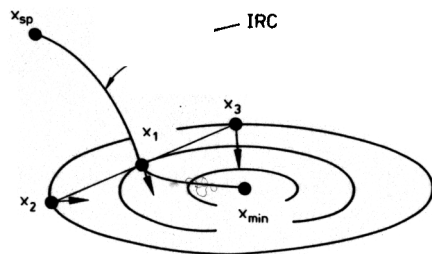


Fig. 2. IRC and its neighborhood (the arrows illustrate length and directions of $-\nabla U$)

strong anharmonic behavior of that section, the conclusions concerning back driving forces connected with this picture are in general wrong. If in a dynamical model by the action of the moment of inertia or, in the computation of the IRC itself by numerical effects we leave the IRC path, say to x_2 or x_3 , then these points will be new starting points for other solutions of the gradient system and no "forces" bring back the system to the IRC. The come-back effect (such a curve is called "stable" by Mezey [60]) does only exist, if the IRC corresponds additionally to a valley path. Furthermore, only in this case it makes a sense to deal with vibrational states orthogonal to the path.

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